

## HYDRODYNAMIC STABILITY THEORY

### Answers to problem sheet 2. Further linear stability.

**Q1.** Given the basic state  $\mathbf{u}_B = (1 - y^2, 0)$ , the NS equations reduce to

$$\partial_x p = \frac{1}{Re} \partial_y^2 u = -\frac{2}{Re}. \quad (1)$$

(Check this by considering each term of each equation in turn.) Integrating gives

$$p_B = -\frac{2x}{Re} + \text{constant}. \quad (2)$$

We now linearise about the basic state. Setting  $\mathbf{u} = \mathbf{u}_B + \delta \tilde{\mathbf{u}}$  gives at  $O(\delta)$

$$\begin{aligned} \partial_t \tilde{u} + u_B \partial_x \tilde{u} + \tilde{v} u'_B &= -\partial_x \tilde{p} + \frac{1}{Re} \nabla^2 \tilde{u}, \\ \partial_t \tilde{v} + u_B \partial_x \tilde{v} + 0 &= -\partial_y \tilde{p} + \frac{1}{Re} \nabla^2 \tilde{v}, \\ \partial_x \tilde{u} + \partial_y \tilde{v} &= 0. \end{aligned} \quad (3)$$

With a normal mode form  $\tilde{u} = \hat{u}(y) \exp[i\alpha(x - ct)]$ , for example, we get

$$\begin{aligned} i\alpha(u_B - c)\hat{u} + \hat{v}u'_B &= -\hat{p}i\alpha + \frac{1}{Re}(D^2 - \alpha^2)\hat{u}, \\ i\alpha(u_B - c)\hat{v} &= -\hat{p}' + \frac{1}{Re}(D^2 - \alpha^2)\hat{v}, \\ i\alpha\hat{u} + D\hat{v} &= 0. \end{aligned} \quad (4)$$

**Q2.** Given  $\mathbf{u} = (0, V(r), 0)$  and  $p = P(r)$ , continuity and the axial equation are trivially satisfied. (Again, check this by examining each term in turn.) The radial equation gives

$$-\frac{V^2}{r} = -P'. \quad (5)$$

The azimuthal equation gives

$$\begin{aligned} \frac{1}{Re} \left[ V'' + \frac{1}{r} V' - \frac{V}{r^2} \right] &= 0, \\ \frac{d^2 V}{dr^2} + \frac{1}{r} \frac{dV}{dr} - \frac{V}{r^2} &= 0, \\ \frac{d}{dr} \left( \frac{d}{dr} + \frac{1}{r} \right) V &= 0. \end{aligned} \quad (6)$$

We now use this to deduce possible forms for  $V(r)$ . Integrating once gives

$$\left(\frac{d}{dr} + \frac{1}{r}\right)V = A \quad \text{with} \quad A = \text{const.} \quad (7)$$

Therefore

$$\frac{1}{r} \frac{d}{dr} (rV) = A. \quad (8)$$

Multiplying across by  $r$  and integrating again gives

$$rV = \frac{Ar^2}{2} + B \quad \text{with} \quad B = \text{const.} \quad (9)$$

So finally we have

$$V = \frac{Ar}{2} + \frac{B}{r}. \quad (10)$$

The perturbation equations are

$$\begin{aligned} \partial_r \tilde{u} + \frac{\tilde{u}}{r} + \frac{1}{r} \partial_\theta \tilde{v} + \partial_z \tilde{w} &= 0, \\ \partial_t \tilde{u} + \frac{V}{r} \partial_\theta \tilde{u} - \frac{2V\tilde{v}}{r} &= -\partial_r \tilde{p} + \frac{1}{Re} \left[ \Delta \tilde{u} - \frac{\tilde{u}}{r^2} - \frac{2}{r^2} \partial_\theta \tilde{v} \right] \\ \partial_t \tilde{v} + \tilde{u} V' + \frac{V}{r} \partial_\theta \tilde{v} + \frac{\tilde{u}V}{r} &= -\frac{1}{r} \partial_\theta \tilde{p} + \frac{1}{Re} \left[ \Delta \tilde{v} - \frac{\tilde{v}}{r^2} + \frac{2\partial_\theta \tilde{u}}{r^2} \right] \\ \partial_t \tilde{w} + \frac{V}{r} \partial_\theta \tilde{w} &= -\partial_z \tilde{p} + \frac{1}{Re} [\Delta \tilde{w}]. \end{aligned} \quad (11)$$

**Q3.** The given form of solution is

$$f(x, y) = 2 \cos(lx\sqrt{3}) \cos(ly) + \cos(2ly). \quad (12)$$

The partial derivatives on the left hand side (LHS) are thus

$$\begin{aligned} \partial_x^2 f &= -6l^2 \cos(lx\sqrt{3}) \cos(ly), \\ \partial_y^2 f &= -2l^2 \cos(lx\sqrt{3}) \cos(ly) - 4l^2 \cos(2ly). \end{aligned} \quad (13)$$

So the LHS is

$$\begin{aligned} \partial_x^2 f + \partial_y^2 f &= -8l^2 \cos(lx\sqrt{3}) \cos(ly) - 4l^2 \cos(2ly) \\ &= -4l^2 \left[ 2 \cos(lx\sqrt{3}) \cos(ly) + \cos(2ly) \right] \\ &= -4l^2 f(x, y). \end{aligned} \quad (14)$$

Therefore we require  $-4l^2 = -a^2$  and so  $a = 2l$ .

Using the relation  $2 \cos(A) \cos(B) = \cos(A - B) + \cos(A + B)$  in (12), we get

$$\begin{aligned} f(x, y) &= \cos(l(x\sqrt{3} - y)) + \cos(l(x\sqrt{3} + y)) + \cos(2ly) \\ &= \cos\left(\frac{a}{2}(x\sqrt{3} - y)\right) + \cos\left(\frac{a}{2}(x\sqrt{3} + y)\right) + \cos(ay). \end{aligned} \quad (15)$$

Therefore we have

$$\begin{aligned}
 f\left(x + \frac{4m\pi}{a\sqrt{3}}, y + \frac{4n\pi}{a}\right) &= \cos\left(\frac{a}{2}(x\sqrt{3} - y) + 2m\pi - 2n\pi\right) \\
 &\quad + \cos\left(\frac{a}{2}(x\sqrt{3} + y) + 2m\pi + 2n\pi\right) \\
 &\quad + \cos(ay + 4n\pi) \\
 &= f(x, y).
 \end{aligned} \tag{16}$$

The last step follows because  $2(m - n)\pi$ ,  $2(m + n)\pi$  and  $4n\pi$  additions leave the cosine function unchanged.

We now proceed to the polar transformation. Writing  $f(x, t) = F(r, \theta)$  with  $x = r \cos \theta$ ,  $y = r \sin \theta$  in (12) we get

$$F = \cos\left(\frac{a}{2}r(\sqrt{3} \cos \theta - \sin \theta)\right) + \cos\left(\frac{a}{2}r(\sqrt{3} \cos \theta + \sin \theta)\right) + \cos(ar \sin \theta). \tag{17}$$

Consider in turn the behaviour of the arguments of each of the three cosine terms in this expression under the transformation  $\theta \rightarrow \theta + \pi/3$ .

- The first argument contains

$$\begin{aligned}
 \sqrt{3} \cos(\theta + \pi/3) - \sin(\theta + \pi/3) &= \sqrt{3} \left\{ \frac{1}{2} \cos \theta - \frac{\sqrt{3}}{2} \sin \theta \right\} - \left\{ \frac{1}{2} \sin \theta + \frac{\sqrt{3}}{2} \cos \theta \right\} \\
 &= -2 \sin \theta.
 \end{aligned} \tag{18}$$

- Similarly the second argument contains

$$\sqrt{3} \cos(\theta + \pi/3) + \sin(\theta + \pi/3) = \sqrt{3} \cos \theta - \sin \theta. \tag{19}$$

- Finally the third argument contains

$$\sin(\theta + \pi/3) = \frac{1}{2} \sin \theta + \frac{\sqrt{3}}{2} \cos \theta. \tag{20}$$

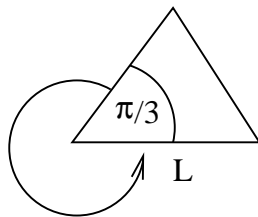
Putting all these together we get

$$\begin{aligned}
 F(r, \theta + \pi/3) &= \cos(-ar \sin \theta) + \cos\left(\frac{ar}{2}(\sqrt{3} \cos \theta - \sin \theta)\right) + \cos\left(\frac{ar}{2}(\sqrt{3} \cos \theta + \sin \theta)\right) \\
 &= F(r, \theta),
 \end{aligned} \tag{21}$$

since  $\cos(-\theta) = \cos(\theta)$ .

P.T.O.

So we have symmetry for rotations by  $\pi/3$ .



Furthermore, since

$$f\left(x + \frac{4m\pi}{a\sqrt{3}}, y\right) = f(x, y) \quad (22)$$

the “wavelength”  $L = 4\pi/a\sqrt{3}$ .